

## Exercise 2

### Question:

The key element of a Xerox machine is a charged drum which is coated with photoconductive material. A simplified sketch of this element is given in fig. 1. Calculate the electric field between the two electrodes.

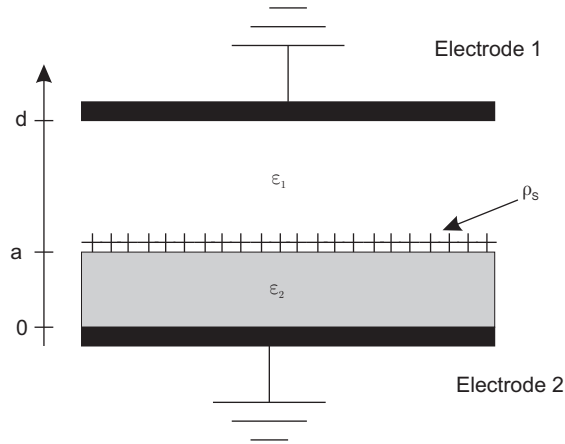


Figure 1: Simplified picture of a xerox machine drum coated with photoconductor. The photoconductor carries a positive charge  $\rho_s$ .

### Solution:

Firstly the domain is split into two half domains at the surface of the photoconductor. Then the governing differential equation is solved in each domain separately:

$$\Delta\Phi = 0 \quad (1)$$

$$\Phi_1(x) = A_1x + B_1 \quad (2)$$

$$\Phi_2(x) = A_2x + B_2 \quad (3)$$

Subsequently, the constants are determined by applying the boundary conditions.

$$\Phi_1(d) = 0 = A_1d + B_1 \quad (4)$$

$$\Phi_2(0) = 0 \rightarrow B_2 = 0 \quad (5)$$

$$\Phi_1(a) = \Phi_2(a) \quad (6)$$

$$D_{1n} - D_{2n} = \rho_s \rightarrow -\varepsilon_1 \frac{d\Phi_1}{dx} + \varepsilon_2 \frac{d\Phi_2}{dx} = \rho_s \quad (7)$$

The constants yield to:

$$A_1 = \frac{-\rho_s}{\varepsilon_1 \left(1 + \frac{\varepsilon_2 d}{\varepsilon_1 a} - \frac{\varepsilon_2}{\varepsilon_1}\right)} \quad (8)$$

$$B_1 = \frac{\rho_s d}{\varepsilon_1 \left(1 + \frac{\varepsilon_2 d}{\varepsilon_1 a} - \frac{\varepsilon_2}{\varepsilon_1}\right)} \quad (9)$$

$$A_2 = \frac{\rho_s \left(\frac{d}{a} - 1\right)}{\varepsilon_1 \left(1 + \frac{\varepsilon_2 d}{\varepsilon_1 a} - \frac{\varepsilon_2}{\varepsilon_1}\right)} \quad (10)$$

$$B_2 = 0 \quad (11)$$

The electric field is determined as follows:

$$E_1 = -A_1 \quad (12)$$

$$E_2 = -A_2 \quad (13)$$

Further detailed information is elaborated in M. N. O. Sadiku, "Elements of Electromagnetics", Saunders College Publishing, Orlando, USA, 1989.

### Exercise 3

Determine a general potential function  $V$  for the region inside the rectangular trough of infinite length in  $z$ -direction whose cross section is shown in fig. 2. Furthermore calculate a specific potential  $V_m$  at  $x = \frac{a}{2}$ ,  $y = \frac{3a}{4}$  for  $V_0 = 100$  V and  $b = 2a$ .

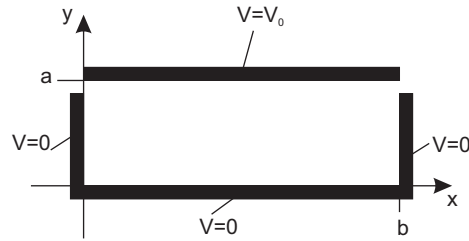


Figure 2: Boundary conditions of the electrical trough.

#### Solution:

Since the trough has an infinite length in the  $z$ -direction the potential  $V$  only depends on  $x$  and  $y$ .

$$\Delta V = \frac{d^2 V}{dx^2} + \frac{d^2 V}{dy^2} = 0 \quad (14)$$

In order to solve this differential equation for the imposed boundary conditions the separation of variables approach is utilised.

$$V(x, y) = X(x)Y(y) \quad (15)$$

Thus the differential equation yields to:

$$X''Y + Y''X = 0 \quad (16)$$

Now we separate to

$$-\frac{X''}{X} = \frac{Y''}{Y} = \lambda. \quad (17)$$

As a separation constant a positive real number  $\lambda = \beta^2$  is chosen. Other separation constants may not satisfy the given boundary conditions. Thus we solve the differential equation to

$$X'' + \beta^2 X = 0 \quad (18)$$

$$X(x) = g_0 \cos(\beta x) + g_1 \sin(\beta x) \quad (19)$$

Imposing the boundary conditions yields:

$$X(x=0) = 0 \rightarrow g_0 = 0 \quad (20)$$

$$X(x=b) = 0 \rightarrow \beta = \frac{n\pi}{b} \quad (21)$$

Thus  $X_n(x) = g_n \sin(\frac{n\pi x}{b})$  is the solution for  $X$ . Similarly we solve

$$Y'' - \beta^2 Y = 0 \quad (22)$$

$$Y(y) = h_0 \cosh(\beta y) + h_1 \sinh(\beta y) \quad (23)$$

Imposing the boundary conditions yields:

$$Y(y=0) = 0 \rightarrow h_0 = 0 \quad (24)$$

Thus  $Y_n(x) = h_n \sinh(\frac{n\pi y}{b})$  is the solution for  $Y$ . The solution in  $x$  and  $y$  is

$$V(x, y) = \sum_n g_n h_n \sinh(\frac{n\pi y}{b}) \sin(\frac{n\pi x}{b}). \quad (25)$$

Now the constant  $c_n = g_n h_n$  has to be determined by the boundary condition

$$V(x, y=a) = V_0 = \sum_n c_n \sinh(\frac{n\pi a}{b}) \sin(\frac{n\pi x}{b}). \quad (26)$$

Integrating with test functions  $\sin(\frac{m\pi x}{b})$  over  $0 < x < b$

$$\int_0^b V_0 \sin(\frac{m\pi x}{b}) dx = \sum_n c_n \sinh(\frac{n\pi a}{b}) \int_0^b \sin(\frac{m\pi x}{b}) \sin(\frac{n\pi x}{b}) dx. \quad (27)$$

With the orthogonality property of the sin function we get

$$\frac{V_0 b}{n\pi} (1 - \cos n\pi) = c_n \sinh(\frac{n\pi a}{b}) \frac{b}{2} \quad (28)$$

This is for  $n$  even

$$c_n = 0. \quad (29)$$

For  $n$  odd we obtain

$$c_n = \frac{4V_0}{n\pi \sinh\left(\frac{n\pi a}{b}\right)}. \quad (30)$$

The complete solution is

$$V(x, y) = \frac{4V_0}{\pi} \sum_i \frac{\sin\left(\frac{(2i-1)\pi x}{b}\right) \sinh\left(\frac{(2i-1)\pi y}{b}\right)}{n \sinh\left(\frac{(2i-1)\pi a}{b}\right)}. \quad (31)$$

Finally the specific potential  $V_m$  is determined to  $V_m = 68.14$  V for a potential  $V_0 = 100$  V,  $b = 2a$  at  $x = \frac{a}{2}$ ,  $y = \frac{3a}{4}$ .

## Exercise 4

An electromagnetic wave described by its electric field as

$$\vec{E} = 4 \sin(2\pi 10^7 t - 0.8x) \vec{e}_z \quad (32)$$

is propagated in a linear, nonmagnetic, homogeneous and isotropic medium. Determine the relative permittivity  $\epsilon_r$  and the wave impedance  $Z = \sqrt{\frac{\mu}{\epsilon}}$  of the medium. Furthermore calculate the time averaged power density carried by the wave as well as the total power crossing an area of  $100 \text{ cm}^2$  in the plane  $2x + y = 5$ .

### Solution:

Extracting from the given electric field

$$\vec{E} = 4 \sin(2\pi 10^7 t - 0.8x) \vec{e}_z = E_0 \sin(\omega t - \beta x) \vec{e}_z \quad (33)$$

equation the following parameter:

$$\begin{aligned} \omega &= 2\pi 10^7 \\ \beta &= 0.8 \\ \mu &= \mu_0 \\ \epsilon &= \epsilon_r \epsilon_0 \end{aligned}$$

Starting with

$$\begin{aligned} \beta &= \omega \sqrt{\mu_0 \epsilon_r \epsilon_0} \\ \epsilon_r &= \frac{\beta^2}{\omega^2 \epsilon_0 \mu_0} \end{aligned} \quad (34)$$

the relative permittivity  $\epsilon_r$  is determined to  $\epsilon_r = \frac{12}{\pi} = 14.59$  and the wave impedance respectively by the following relation:

$$Z = \sqrt{\frac{\mu_0}{\epsilon}} = 10\pi^2 = 98.7\Omega. \quad (35)$$

The Pointing vector is defined as

$$\vec{S} = \vec{E} \times \vec{H}. \quad (36)$$

The magnetic field strength has to be derived first. We start from Faraday's law

$$\begin{aligned} \nabla \times \vec{E} &= -\mu_0 \frac{d\vec{H}}{dt} \\ \vec{e}_y \frac{dE}{dx} &= E_0 \beta \cos(\omega t - \beta x) \vec{e}_y \\ \vec{H} &= E_0 \frac{\beta}{\mu_0} \vec{e}_y \int \cos(\omega t - \beta x) dt \\ \vec{H} &= -E_0 \frac{\beta}{\omega \mu_0} \sin(\omega t - \beta x) \vec{e}_y \end{aligned}$$

and utilising the relation

$$\frac{\beta}{\omega \mu_0} = \frac{\omega \sqrt{\mu_0 \epsilon}}{\omega \mu_0} = \sqrt{\frac{\epsilon}{\mu_0}} = \frac{1}{Z}$$

the magnetic fields yield to

$$\vec{H} = -\frac{E_0}{Z} \sin(\omega t - \beta x) \vec{e}_y = \frac{E}{Z} \vec{e}_y. \quad (37)$$

Consequently the pointing vector yield to

$$\vec{S} = \frac{E_0^2}{Z} \sin^2(\omega t - \beta x) \vec{e}_x. \quad (38)$$

and finally the time averaged pointing vector can be described as

$$\vec{S}_{ave} = \vec{e}_x \frac{1}{T} \int_0^T \frac{E_0^2}{Z} \sin^2(\omega t - \beta x) dt = 81.05 \vec{e}_x \frac{\text{mW}}{\text{m}^2}, \quad (39)$$

where  $T = \frac{2\pi}{\omega}$  describes the time period. The total power crossing the plane  $2x + y = 5$  with an angle  $\varphi = \arctan(\frac{1}{2})$  with respect to its normal vector  $\vec{n} = (2, 1, 0)^T / \sqrt{5}$ , finally yields to:

$$\begin{aligned} P_{avg} &= \vec{S}_{ave} \cdot \vec{A} = \vec{S}_{ave} \cdot \vec{n} A \\ &= \vec{e}_x \cdot \vec{n} A S_{ave} \\ &= \cos(\arctan(\frac{1}{2})) A S_{ave} = 724.5 \mu\text{W}. \end{aligned} \quad (40)$$